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Quadrupole Moment for High Spin Rotational States

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The asymmetric rotor model is proposed to explain the behaviors of high spin states,^{1,2)} where the pairing force has no significant effect and the asymmetric rotor is more stable than the symmetric one. Recently there comes a lots of experiments about the yrast line,³⁾ back bending curve,⁴⁾ and the rotational spectra of unique parity states in odd *A* nucleus.⁵⁾ All these phenomena, especially the fast E2-transition rate, which seems to be essential peculiarity in yrast region, are interpreted in terms of this model. The rotational Hamiltonian is diagonalized by Holstein-Primakoff transformation, and a physical quantity is expanded in powers of 1/*I* for the spin *I* state. The quadrupole moment is given by

$$Q = \frac{z}{m} \frac{1}{(I+1)(2I+3)} \sum_{n=0}^{2I} \left[A \{ 3(I-n)^2 - I(I+1) \} G_{nk}^I(x) + B\sqrt{(2I-n)(2I-n+1)(n+1)(n+2)} \times G_{n+2,k}^I(x) \right] G_{nk}^I(x),$$

where $A = \mathscr{I}_x + \mathscr{I}_y - 2\mathscr{I}_z$, and $B = 3(\mathscr{I}_x - \mathscr{I}_y)$. In the above \mathscr{I} 's are the moments of inertia, Z is charge and m mass. Here the transformation coefficients $G_{kl}^I(x)$ are explicitly given by⁶

$$G_{kl}^{I}(x) = \frac{(-)^{[k/2]}}{\left(\frac{k+l}{2}\right)!} \left(\frac{k!l!}{2^{k+l}\cosh^{k+l+1}(2x)}\right)^{1/2} \\ \times \sum_{p=0}^{\min([k/2],[l/2])} (-4)^{p}\sinh^{[k/2]+[l/2]-2p}(2x) \\ \times \left(\left[\frac{k}{2}\right] + \left[\frac{l}{2}\right] - 2p \\ \left[\frac{k}{2}\right] - p\right) \times \left(\frac{\binom{k+l}{2}}{2p}, \text{ for } k, l \text{ even,} \\ 2\binom{\binom{k+l}{2}}{2p+1}, \text{ for } k, l \text{ odd} \right)$$

A set of parameters which is consistent with the results of ref. 4| is| given by $\mathscr{I}_x = 125$, $\mathscr{I}_y = 25$ and $\mathscr{I}_z =$ 0.9 (in the unit of MeV⁻¹). Then, $Q^{\text{asym}}(20+)/Q^{\text{sym}}(20+) = 2.7$. The calculated yrast line is shown in the figure.



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