

Table IV
Magnetic Moments of Shell-Model States in the Lead Region

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In this table, nuclear g factors of several shell-model states in the Pb region are listed.

- (i) Since there is inconsistency in the corrections for the diamagnetism and the Knight shift between published data by different authors, g factors before the corrections are listed; $g_{\text{exp}}(\text{uncorr})$. Then the following corrections are applied to obtain the "Adopted g factor". For measurements in non-magnetic insulator hosts

$$\Delta H/H = -1.4 \pm 0.4\% \quad (\text{diamagnetic correction}).^{1)}$$

For metallic hosts

$$\Delta H/H = 0.0 \pm 1.0\%$$

$$\begin{aligned} & (\text{diamagnetic correction } (\approx -1.4\%) \\ & + \text{ Knight-shift correction } (\approx 1.5\%)).^{2)} \end{aligned}$$

- (ii) In the case where two or more experimental values of the g factor of the same state were reported, a simple average (without weights) of them after the corrections is listed in the column of the "Adopted g factor".
- (iii) Values reported at this conference are indicated by underlines.

(1) Single-particle states

Nucleus	J^π	E^* (keV)	$T_{1/2}$	g_{exp} (uncorr) ^{a)}	Ref.	Adopted g factor	Main configuration	g_{Schmidt}
²⁰⁷ Pb	5/2 ⁻	570	129ps	0.288(28)	3	0.29(3)	$\nu f_{5/2}^{-1}$	0.547
				0.26(2)	3			
				0.316(15)	4			
				1.1673	5	1.184(5) ^{b)}	$\nu p_{1/2}^{-1}$	1.275
²⁰⁷ Pb	1/2 ⁻	gnd	stable				$\pi h_{9/2}$	0.583
²⁰⁹ Bi	9/2 ⁻	gnd	$> 2 \times 10^{18}$ y	0.8975	5	0.910(4) ^{b)}	$\pi s_{1/2}^{-1}$	5.586
²⁰⁵ Tl	1/2 ⁺	gnd	stable	3.223	5	3.269(13) ^{b)}	$(\nu p_{1/2}^{-2})_{0^+}$	0.294
²⁰⁵ Pb	13/2 ⁺	1014	5.55ms	-0.150(4)	6	-0.150(6)	$(\nu p_{1/2}^{-2})_{0^+}$	1.655
²¹¹ Bi	7/2 ⁻	405	318ps	1.26(19)	3	1.26(19)	$\pi f_{7/2}$	
²⁰⁷ Po	13/2 ⁺	1114	47 μ s	-0.1427(15)	7	-0.143(2)	$\nu i_{13/2}^{-1}$	0.294
²¹¹ Po	15/2 ⁻	1065	16ns	-0.05(2)	8	-0.05(2)	$\nu j_{15/2}$	-0.255 ^{c)}
							+ (3 ⁻ $\otimes$ $\nu g_{9/2}(\pi h_{9/2})_{15/2} - \otimes (\pi h_{9/2}^2)_{0^+}$	

a) Experimental values without any corrections for the Knight shift and chemical shift.

b) Corrected for the chemical shift of $-1.4 \pm 0.4\%$ ¹⁾

c) Schmidt value for the $\nu j_{15/2}$ state.

(2) Many-particle states, $(j_1 \otimes j_2)J$, with $j_1 \neq j_2$.

Nucleus	J^π	E^* (keV)	$T_{1/2}$	g_{exp} (reported)		Ref.	Adopted g factor	Main configuration	$g_{\text{cal}}^{\text{d)}$ (add.)	g_{Schmidt}
				uncorrected	corrected					
^{203}Tl	7^+	950	0.56ms	0.128(6)		9	0.128(6)	$\pi s_{1/2} \nu i_{13/2}^{-1} \otimes (^{204}\text{Pb})_0^+$	0.09 ^{e)}	0.126
^{206}Pb	6^-	2385	29ps	0.13(7)		10	0.13(7)	$\nu p_{1/2} \nu i_{13/2}^{-1}$	-0.25 ^{f)}	-0.406
^{208}Pb	7^-	2200	123 μ s	-0.0217(2)	-0.0217(4)	11	-0.0217(4)	$\nu p_{1/2} \nu i_{13/2}^{-1}$	-0.06 ^{f)}	-0.182
^{208}Pb	5^-	3198	297ps	-0.013(27)		12	0.02(1)	$\nu g_{9/2} \nu p_{1/2}^{-1}$	(-0.15) ^{g)}	-0.255
				0.057(8)		13				
				0.021(7)		14				
^{206}Bi	6^+	gnd	6.24d		0.76(1)	3	0.76(1)	$\pi h_{9/2} \nu f_{5/2}^{-1} \otimes (\nu p_{1/2}^{-1})_0^+$	0.72	0.572
^{207}Bi	$21/2^+$	2102	182 μ s	0.325(3)	0.325(6)	11	0.325(6)	$(\nu p_{1/2}^{-1} \nu i_{13/2})_7^+ \otimes \pi h_{9/2}$	0.32	0.101
^{210}Bi	5^-	439	38ns ^{a)}	(0.309(9))	0.306(9)	15	0.309(9)	$\pi h_{9/2} \nu g_{9/2}$		0.079
^{210}Bi	7^-	433	58ns ^{a)}	(0.305(7))	0.302(7)	15	0.305(8)	$\pi h_{9/2} \nu g_{9/2}$		0.079
^{209}Po	$17/2^-$	1473	98ns ^{b)}	0.88(5)		16	0.890(15)	$(\pi h_{9/2})_8^+ \otimes \nu p_{1/2}^{-1}$	0.926	0.624
					0.897(15)	17				
^{210}Po	11^-	2849	24ns	1.107(16)		18	1.107(19)	$\pi h_{9/2} \pi i_{13/2}$		1.038
^{210}Po	13^-	4392	93ns		0.546(12)	19	0.546(12)	$(\pi h_{9/2})_8^+ \otimes 5^-(\text{in } ^{208}\text{Pb})$	0.56	(0.339)
^{210}At	15^-	2549	0.75 μ s ^{c)}	1.044(6)		20	1.04(1)	$[(\pi h_{9/2})_8^+ \otimes \pi i_{13/2}]_{29/2} + 1.07^{\text{h)}$		0.939
								$\otimes \nu p_{1/2}^{-1}$		
^{211}At	$29/2^+$	2641	70ns		1.03(4)	21	1.03(4)	$(\pi h_{9/2})_8^+ \otimes \pi i_{13/2}$	1.07 ^{h)}	0.928

- a) Ref. 22. b) Ref. 23. c) Ref. 24.
d) Calculated from the experimental g factors of the single-particle states with the additivity relation of the g factor.
e) $g(\pi s_{1/2}^{-1}; ^{203}\text{Tl})$ and $g(\nu i_{13/2}^{-1}; ^{205}\text{Pb})$ are used.
f) $g(\nu i_{13/2}^{-1}; ^{206}\text{Pb})$ is used for the g factor of the $\nu i_{13/2}^{-1}$ state.
g) $g(\nu g_{9/2}) = -0.30(1)$ deduced from the g factors of the 5^- and 7^- states of ^{210}Bi is used.
h) $g(\pi i_{13/2}) = 1.268(36)^{25)}$ deduced from the g factor of the 11^- state of ^{210}Po is used.

(3) Proton- $h_{9/2}$ n-particle states.

Nucleus	J^π	E^* (keV)	$T_{1/2}$	g_{exp} (reported)		Ref.	Adopted g factor	Main configuration
				uncorrected	corrected			
^{209}Bi	$9/2^-$	gnd	$> 2 \times 10^{18}\text{y}$	(See Table (1))			0.910(4)	$\pi h_{9/2}$
^{210}Po	8^+	1557	110ns	0.925(11)		26	0.910(10)	$\pi h_{9/2}^2$
				0.904(14)		18		
					0.909(15)	27		
					<u>0.901(13)</u>	17		
^{210}Po	6^+	1473	41ns^{a}	0.93(2)		27	0.93(2)	$\pi h_{9/2}^2$
^{211}At	$21/2^-$	1417	50ns	<u>0.897(16)</u>		28	0.901(20)	$\pi h_{9/2}^3$
^{212}Rn	8^+	≈ 1690	$1.0\mu\text{s}$	0.89(3)		29	0.89(3)	$\pi h_{9/2}^4$
^{204}Po	8^+	≈ 1650	140ns^{b}	0.905(40)		17	0.97(8)	$\pi h_{9/2}^2 \otimes (vp_{1/2}^{-2}vf_{5/2}^{-4})_0^+$
				<u>1.04 (8)</u>		30		
^{206}Po	8^+	≈ 1590	212ns^{c}	0.926(16)		31	0.93(3)	$\pi h_{9/2}^2 \otimes (vp_{1/2}^{-2}vf_{5/2}^{-2})_0^+$
				0.905(18)		17		
				<u>0.95 (4)</u>		30		
^{208}Po	8^+	≈ 1530	380ns	0.915(6)		32	0.913(10)	$\pi h_{9/2}^2 \otimes (vp_{1/2}^{-2})_0^+$
						17		
^{209}Po	$17/2^-$	1473	98ns	<u>0.911(10)</u>			0.890(15)	$(\pi h_{9/2}^2)_8^+ \otimes vp_{1/2}^{-1}$
				(See Table (2))				

a) Ref. 33. b) Ref. 34. c) Ref. 35.

(4) Neutron- $i_{13/2}$ n-particle states

Nucleus	J^π	E^* (keV)	$T_{1/2}$	g_{exp} (reported)		Ref.	Adopted g factor	Main configuration
				uncorrected	corrected			
^{205}Pb	$13/2^+$	1014	5.55ms	-0.150(4)	-0.150(6)	6	-0.150(6)	$vi_{13/2}^{-1} \otimes (vp_{1/2}^{-2})_0^+$
^{206}Pb	12^+	4027	200ns	-0.1547(35)	-0.155(4)	36	-0.155(4)	$vi_{13/2}^{-2}$
^{207}Po	$13/2^+$	1114	$47\mu\text{s}$	-0.1427(15)	-0.143(2)	7	-0.143(2)	$vi_{13/2}^{-1} \otimes (vp_{1/2}^{-2}\pi h_{9/2}^2)_0^+$

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