## **Dielectric Properties of Dirty Ferroelectric SBN**

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Diffuse phase transition of a rare earth iondoped SBN40 ( $Ba_{0.4}Sr_{0.6}Nb_2O_6$ ;  $Pr^{3+}$  or  $Nd^{3+}$ ) has been carefully examined by the measurements of the spontaneous polarization  $P_{\rm s}$ , the dielectric constant  $\varepsilon$  and X-ray diffraction. P<sub>s</sub> decreases gradually with increasing temperature and  $\varepsilon$  has a broad maximum near the transition temperature  $T_c$ . With increasing rare earth ion concentration M, 1) the broadening of  $\varepsilon$  increases, 2) the temperature dependence of  $P_s$  becomes more diffuse, 3)  $T_c$ and  $T_{p}$  (the inflexion point of  $P_{s}$ ) decrease in proportion to  $M^{-1}$ . In order to discuss the temperature and M dependence of  $P_{s}$  and  $\varepsilon$ , we adopt the order parameter expansion of the free energy including odd powers as well as even powers of the polarization and the expansion of the reciprocal susceptibility in terms of  $T - T_c$ which was introduced by L. Benguigui:1)

$$\chi^{-1} = \chi_0^{-1} + B_2 (T - T_c)^2 + B_3 (T - T_c)^3.$$
(1)

 $B_2$  and  $B_3$  are the expansion coefficients, and were determined by the method of least square. It turned out that  $B_2/B_3$  is proportional to  $M^{-1}$ . The calculated values of  $\chi^{-1}$  are in good agreement with the experimental results in the temperature range of  $T_{\rm c} \pm 5^{\circ}$ C for pure SBN40, and  $T_{\rm c} \pm 15^{\circ}$ C for the samples with the rare earth ion. However, except the above temperature regions,  $\chi^{-1}$  is a linear function of temperature. The ratio of the slopes of the  $\chi^{-1}$  vs T curve at low and high temperature regions is about 2:1 in our experiments. Figure 1 shows the temperature dependence of  $P_s$ . Marks show the calculated values and the lines are the experimental results. The calculated values of  $P_s$  agree with the experimental results in the temperature range of  $T_{\rm c} \pm 20^{\circ}$ C for the samples containing the rare earth ion. Furthermore, by using the expressions for  $B_2$  and  $B_3$ ,<sup>1)</sup> we can write the

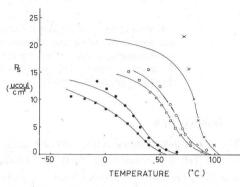


Fig. 1. Temperature dependence of the spontaneous polarization  $P_s$  and the calculated values of  $P_s$  for various samples. The lines show the experimental results.×: pure SBN40,  $\bigcirc$ : SBN40+Pr<sup>3+</sup> (0.3 wt%),  $\bigcirc$ : SBN40 +Pr<sup>3+</sup> (1.0wt%),  $\Box$ : SBN40+Nd<sup>3+</sup> (0.3 wt%),  $\blacksquare$ : SBN40+Nd<sup>3+</sup> (1.0 wt%)

inflexion point as

$$T_{\rm p} = T_{\rm c} + B_2 / 6B_3. \tag{2}$$

It is clear that the calculated values of  $T_p$  given by eq. (2) agree with the experimental values within 5% error in all samples. As mentioned above, in the case of the rare earth ion-doped SBN40, temperature dependence of the dielectric properties is well explained by the order parameter expansion of free energy and the expansion of  $\chi^{-1}$  in terms of  $T - T_c$ . Furthermore, X-ray experiments were carried out to clarify the microscopic mechanism of the diffuse phase transition. The equi-intensity contour of X-ray critical diffuse scattering around (0 0 4) reflection in SBN40 +  $Pr^{3+}$  (0.3wt%) shows a strongly anisotropic distribution. Considering the piezoelectric coupling in  $\chi^{-1}$ , we can explain the strong anisotropy near  $T_c$  in SBN40 + Pr<sup>3+</sup> (0.3wt%).

## Reference

1) L. Benguigi: Solid State Commun. 14 (1974) 669.