## Dielectric Critical Slowing-Down and Some Related Properties in RbH<sub>3</sub>(SeO<sub>3</sub>)<sub>2</sub>

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Dielectric, ultrasonic, thermal and X-ray measurements for  $\text{RbH}_3(\text{SeO}_3)_2$  were made near the ferroelectric transition temperature  $T_c$ .  $\text{RbH}_3(\text{SeO}_3)_2$  exhibits anomalies in the ultrasonic velocities and in the specific heat both at  $T_i = (2.2^\circ\text{C} + T_c)$ , the upper transition temperature and  $T_c$ . X-ray examinations confirmed an incommensurate nature of the intermediate phase below  $T_i$ . The dielectric results obtained above  $T_c$  in a frequency range between 3 MHz and 1000 MHz show critical slowing-down of the relaxation time, which is suggested to be originated from discommensuration of the incommensurate phase.

Because of unusual dielectric properties of rubidium trihydrogen selenite, RbH<sub>3</sub>(SeO<sub>3</sub>)<sub>2</sub>, such as a small peak of the dielectric constant at the ferroelectric transition temperature  $T_{\rm c}$  and extremely small spontaneous polarization below  $T_{c}^{(1)}$  many investigations<sup>2-4</sup> have been made to understand its ferroelectric transition mechanism. Dvorak<sup>2)</sup> and Sanikov et al.<sup>3)</sup> analyzed irreducible representations of RbH<sub>3</sub>(SeO<sub>3</sub>)<sub>2</sub> for the ferroelectric phase transition, and the latter authors discussed a possibility of an intermediate incommensurate phase of  $RbH_3(SeO_3)_2$  just above  $T_c$ . Recently Gladkii et al.5) found a small anomaly in the elastic compliance  $s_{55}^{E}$  at about 2 K above  $T_{c}$  and attributed the anomaly to an incommensurate to the paraelectric (commensurate) phase transition. Very recently Gesi et al.<sup>6)</sup> made neutron diffraction study of the deuterated crystal  $RbD_3(SeO_3)_2$  and reported an incommensurate nature of the intermediate phase.

We have recently measured for  $RbH_3(SeO_3)_2$ 

(1) the ultrasonic velocities of both longitudinal and transverse waves  $v_i$  ( $i=1, 2, \dots, 6$ ) and absorption coefficients  $\alpha_{ii}$  (i=1, 2, 3),

(2) the specific heat,

(3) the complex dielectric constant at frequencies below 1000 MHz, and examined

(4) incommensurate nature by the X-ray diffractometer.

The results of (1) and (3) were shortly reported elswhere.<sup>7,8)</sup>

The results of the ultrasonic velocities of the longitudinal waves  $v_i$  (i=1, 2, 3) and of the transverse waves  $v_i$  (i=4, 5, 6) are shown in Fig. 1. All the ultrasonic velocities  $v_i$ 's except  $v_6$  show anomalies both at ( $2.2^{\circ}C + T_c$ ) and at  $T_c$ . In what follows, the temperature of ( $2.2^{\circ}C + T_c$ ) is denoted as the incommensurate to the paraelectric (commensurate) phase transition tem-



Fig. 1. Sound velocities vs temperature curves. q and  $\xi$  show the wave vector and the polarization vector, respectively.

perature  $T_i$ , because of the appearance of incommensurate nature of the intermediate phase as will be mentioned below. Figure 2 demonstrates the temperature dependence of  $v_3$ and  $\alpha_{33}$  in the very vicinity of  $T_i$ , as an example of the results showing the anomalies in  $v_i$ 's at  $T_c$ and  $T_i$ . It is noted that the anomalous parts of the observed velocities and absorption coefficients,  $\Delta v_i = (v_i^0 - v_i)$  and  $\Delta \alpha_{ii}$  (i=1, 2, 3), all show a tendency to diverge at  $T_i$ . The normal part of the velocity  $v_i^0$  near  $T_i$  is obtained by extrapolating  $v_i$  vs T curve in the high temperature range above  $T_i$  to low temperature.

The more direct evidence for the existence of the intermediate phase was obtained by the measurement of the specific heat  $C_p$ , the result of which is shown in Fig. 3. The total transition entropy  $\Delta S_t$  associated with both the lower (ferroelectric) and upper (incommensurate to paraelectric) phase transitions was estimated to be 0.12 cal/mol · deg.

The temperature dependence of the real and imaginary parts of the complex dielectric constant along the *b* axis,  $\varepsilon'_b$  and  $\varepsilon''_b$ , measured at various frequencies between 3 MHz and 1000 MHz are illustrated in Fig. 4. The results shown in the figure demonstrate the dielectric dispersion observed near  $T_c$ . The real part  $\varepsilon'_b$ measured at each frequency except 3 MHz and



Fig. 2. Sound velocity  $v_3$  and absorption coefficient  $\alpha_{33}$  of the longitudinal wave with the wave vector q//[001] and the polarization vector  $\xi//[001]$  neat  $T_i$ .



Fig. 3. The specific heat as a function of temperature.



Fig. 4. The real and imaginary parts of the complex dielectric constant,  $\varepsilon'_{b}$  and  $\varepsilon''_{b}$ , neat  $T_{c}$ .

1000 MHz indicates a discontinuous change in value at  $T_c$ , indicating that the ferroelectric transition of RbH<sub>3</sub>(SeO<sub>3</sub>)<sub>2</sub> at  $T_c$  is of first order. Examination of the dielectric result by the Cole-Cole relation,  $\varepsilon(\omega) - \varepsilon(\infty) = {\varepsilon(0) - \varepsilon(\infty)}/{1 + (i\omega\tau)^{\beta}}$ , indicates that the value of a measure of the width of the distribution of relaxation times,  $\beta$ , falled into a range  $0.9 < \beta < 1$  with uncertainty  $\Delta\beta = 0.1$ . The relaxation time obtained in the incommensurate phase shows critical slowing-down with a value of  $1.7 \times 10^{-9}$  sec just above  $T_c$ . However, RbH<sub>3</sub>(SeO<sub>3</sub>)<sub>2</sub> can be regarded as of the displacive type by the observation of a

zone-boundary soft phonon mode,<sup>9)</sup> for which the dielectric dispersion is expected to arise at frequencies of the order of lattice vibrations. The relaxation frequencies actually observed in  $RbH_3(SeO_3)_2$  are too low to attribute to those of the lattice vibrations.

According to McMillan,<sup>10)</sup> in the vicinity of  $T_c$  in the incommensurate phase the crystal has a domain-like structure with discommensuration. The discommensuration distance increases as  $T_c$  is approached from above. It is therefore suggested that the critical slowing-down observed in RbH<sub>3</sub>(SeO<sub>3</sub>)<sub>2</sub> near  $T_c$  may be due to the relaxation of polarization contributed from the domain-wall-like movement of discommensuration.

Incommensurate nature of the intermediate phase of  $\text{RbH}_3(\text{SeO}_3)_2$  reported by Gesi *et al.* was confirmed by the X-ray examination. The results of X-ray and neutron<sup>6</sup> diffractions thus

support the explanation for the origin of the observed critical slowing-down.

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