

1.12 Determination of the Spin-Spin Potential in the Optical Model from the Depolarization of Proton Elastic Scattering *)

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For the elastic scattering of nucleons from target nuclei with non-zero spin a spin-spin potential has been suggested ¹⁾ in addition to the well-known real, imaginary and spin-orbit potential terms. The spin-spin interaction is considered to be the sum of a spherical term

$$U_{SS}(r) = -V_{SS} f_O(r) \underline{\sigma} \cdot \underline{I}$$

and a tensor term

$$U_{ST}(r) = -\frac{1}{2} V_{ST} f_T(r) \cdot [3(\underline{\sigma} \hat{r}) - (\underline{\sigma} \cdot \underline{I})]$$

where $\underline{\sigma}$ and $\underline{I}$ denote particle and target spin, respectively, $\hat{r}$ the relative vector and f_O and f_T the appropriate form factors. Theoretical estimates in the framework of different models ²⁻⁴⁾ resulted in spherical spin-spin potentials much smaller than the usual spin orbit potential, whereas the tensor spin-spin potentials should be even smaller. The most sensitive observable for these spin-spin potentials is the depolarization $D(\theta)$, which can be obtained from the analyzing power $A(\theta)$ and the polarization of incoming and outgoing particles.

In this contribution we report on the measurement and a full analysis of $\sigma(\theta)$, $A(\theta)$ and $D(\theta)$ for the elastic proton scattering on ²⁷Al ($E_p=10.35$ MeV and 11 MeV) and on ⁸⁹Y ($E_p=11$ MeV) in the angular range $40^\circ \leq \theta_{Lab} \leq 120^\circ$. Differential cross section and analyzing power were measured in the usual way, whereas the polarization of the scattered protons was obtained in a special double scattering arrangement ⁵⁾ consisting of a sliding seal chamber, a high acceptance QDQ magnetic spectrometer with $\Omega=20$ msr and a high pressure ⁴He-polarimeter ($p=20$ bar). For the elimination of instrumental asymmetries both the particle spin and the polarimeter are turned around 180° in a different time scale. The polarization of the incoming beam was measured on line in the first scattering chamber, where the known analyzing power of ²⁷Al and ⁸⁹Y, respectively, at backward angles was used.

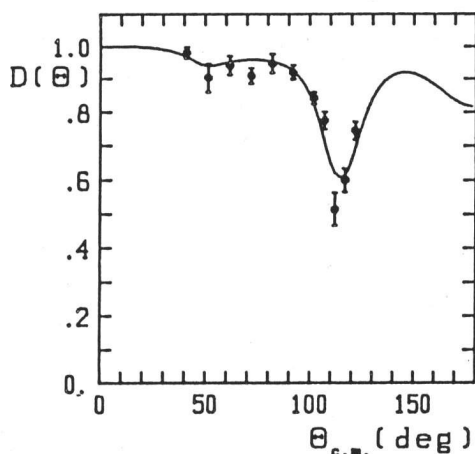


Fig.1 Measured depolarization for ²⁷Al($\vec{p}, \vec{p}_O$) at 11 MeV with a best fit curve ($V_{SS}=1.1$ MeV; $V_{ST}=0.8$ MeV).

As an example the measured depolarization is shown for ²⁷Al($\vec{p}, \vec{p}_O$) at $E_p=11$ MeV compared with a theoretical curve, where besides the spin-spin interaction compound nucleus effects ⁶⁾ and the quadrupole spin-flip ⁷⁾ is taken into account properly. The optical potential without spin-spin terms was obtained from an analysis of $\langle \underline{\sigma} \cdot \underline{A} \rangle$, which depends on the direct scattering only, and $\underline{\sigma} \cdot \underline{DI} = \langle \underline{\sigma} \rangle - \underline{\sigma}^{CE}$, where the com-

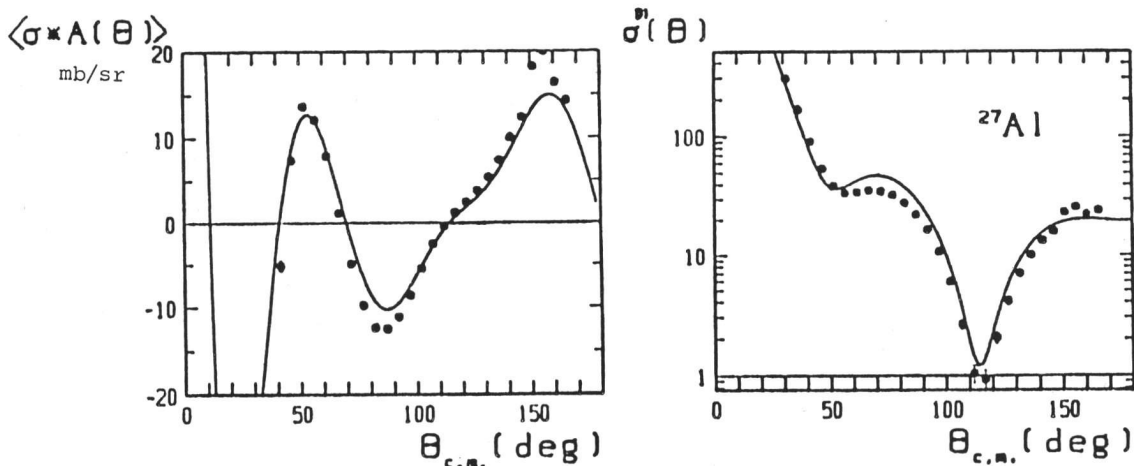


Fig.2 $\langle \sigma \cdot A \rangle$ and σ^{DI} for $^{27}\text{Al}(\vec{p}, p_0)$ at $E_p=11$ MeV compared with optical model curves.

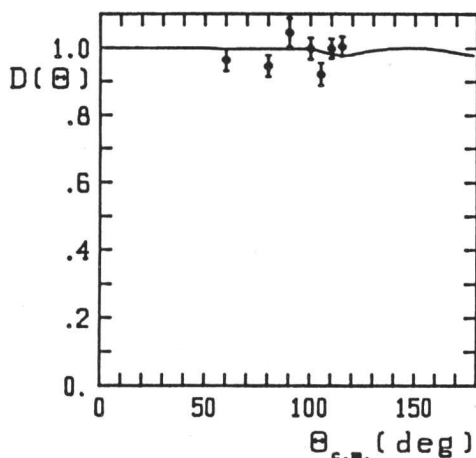


Fig.3 Measured depolarization for $^{89}\text{Y}(\vec{p}, p_0)$ at $E_p=11$ MeV compared with a theoretical curve ($V_{SS}=V_{ST}=1$ MeV).

pound-elastic contribution was calculated within a modified Hauser-Feshbach formalism (see fig.2). For ^{27}Al the results at both energies are consistent with $V_{SS}=(1+0.2)$ MeV and $V_{ST}=1$ MeV. For the $I=1/2$ nucleus ^{89}Y , where the effects in $D(\theta)$ are much smaller, the depolarization data only yield an upper limit of 1 MeV for V_{SS} and V_{ST} (see fig.3).

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