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Competition of Two-Step Sequential Transfer Processes with Direct Process in the Reaction  $^{208}Pb(d,\alpha)^{206}T1(0,g.s.)$  and a Complete Experiment of the Transition

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Reaction mechanism of (d,  $\alpha$ ) and/or ( $\alpha$ , d) reaction has been discussed recently.<sup>1</sup> In addition to the direct one-step  $(d, \alpha)$  process, it has pointed out that there exist two-step processes and the contribution of the D state of the  $\alpha$  particle. In this report we take the case of reaction  $^{208}\text{Pb}(d,\alpha)^{206}\text{T1}(0\overline{g}_{s},(\pi s_{1/2})^{-1}(\nu p_{1/2})^{-1})$  and present the evidence of strong sequential two-step processes in the  $(d, \alpha)$  reaction.

We measured the cross section, vector-  $(iT_{11})$  and tensor analyzing powers  $(T_{20}, T_{21}, T_{22})$  of this reaction at  $E_d = 21.5$  MeV. The data were analyzed in terms of the full finite range (FR) first- and second-order distorted wave Born approximation (DWBA). This reaction is explained in terms of the sequential transfer  $(d,t)(t,\alpha)$ and (d, <sup>3</sup>He)(<sup>3</sup>He,  $\alpha$ ) two-step processes in addition to a (d,  $\alpha$ ) one-step process by taking account of  $\alpha$  particle D-state. The one-step (d,  $\alpha$ ) process is generated by the deuteron-deuteron cluster interaction. The D-state admixture is chosen so that the parameter  $\rho$ , the asymptotic ratio of the S- and D-state wave function is -0.38<sup>2</sup>? The absolute value of the one-step cross section was normalized so that the total theoretical cross sections obtained from a coherent sum of the one- and two-step transition amplitudes reproduce the experimental  $(d, \alpha)$  cross sections. The strength of each of the one-nucleon transfer reactions 3), 4) in the sequential transfer two-step processes is determined by comparing the experimental cross sections of each of the one-nucleon transfer reactions with the corresponding FR DWBA cross sections (Fig. 1). The intensity of the two-step processes is as large as that of the one-step process so that the constructive interference between the one- and two-step transition amplitudes play an essential role in reproducing not only the observed analyzing powers but also the experimental cross section (Fig. 2). The D-state effect is minor.

This reaction has a special spin-and-parity sequence of  $f + d \rightarrow d + 0$ . can express the transition matrix M for this reaction as M = [<00 | M | 10 > <00 | M | 00 > $\langle 00 | M | -10 \rangle \equiv [A B C]$ , where the notations  $\alpha$ ,  $\beta$ ,  $\gamma$  and  $\delta$  in the matrix element  $<\gamma\delta|M|\alpha\beta>$  denote the magnetic quantum numbers of the particles d, <sup>208</sup>Pb,  $\alpha$  and <sup>206</sup>T1 respectively. According to the Bohr's theorem<sup>5</sup>), the transition matrix element C = -A. Therefore, this transition can be described by A and B. So we obtained the = -A. Inererore, this transition can be described by A and B. So we obtained the relations for the analyzing powers, A  $_{yy} = -(T_{20} + \sqrt{6} T_{22}) / \sqrt{2} = 1$ , and  $|iT_{11}|^2 + |T_{20}|^2 / 2 + |T_{21}|^2 + |T_{22}| = 1$ , then we can calculate the T<sub>22</sub> and the absolute value of T<sub>21</sub> only by using measured values of the iT<sub>11</sub> and T<sub>20</sub>. Each of the analyzing powers is described by using the transition matrix elements A and B as  $iT_{11} = \sqrt{6}$   $Im(AB^*)/I_0$ ,  $T_{20} = \sqrt{2}$  ( $|A|^2 - |B|^2)/I_0$ ,  $T_{21} = -\sqrt{6}$  Re(AB\*)/I<sub>0</sub>, and  $T_{22} = -\sqrt{3} |A|^2/I_0 < 0$ , where  $I_0 = 2 |A|^2 + |B|^2 = 3\sigma_0$ , and  $\sigma_0$  is the unpolarized cross section. We put  $A = |A| e^{4\alpha}$ , and  $B = |B| e^{4\beta}$ . Then they become  $A = [\sigma_0(T_{20}/\sqrt{2} + 1)]^{1/2}$ ,  $B = [\sigma_0(1 - \sqrt{2}T_{20})]^{1/2}$ ,  $sin(\alpha - \beta) = \sqrt{3} iT_{11} / \sqrt{2} (1 - T_{20}/\sqrt{2} - T_{20}^2)^{1/2}$  and  $tan(\alpha - \beta) = -iT_{11} / T_{21}$ . Thus all the transition matrix elements A and B are completely determined by  $T_{21}$ . Thus all the transition matrix elements A and B are completely determined by the present experiment.

## References

- 1) W. W. Daehnick et al., Phys.Rev.C23(1981)1906. W. T. Pinkston et al., Nucl.Phys. A383(1982)61. F. D. Santos et al., Phys.Rev.C25(1982)3243. M. A. Nagarajan et al., Phys.Rev.Lett.47(1982)1899. J. R. Comfort et al., ibid.50(1983)1627.

- J. A. Tostevin, Phys.Rev.C28(1983)961.
  P. D. Barnes et al., Phys. Rev. C1(1978)228.
  W. D. Alford and D. G. Burke, Phys. Rev.185(1973)2366.
- 5) A. Bohr, Nucl. Phys. <u>10</u>(1959)486.

