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7.2 Energy Dependence of the j-Effect in Radiative Capture of Polarized Nucleons

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It has recently been pointed out<sup>1)</sup> that polarized nucleon capture data indicate a sensitivity to the  $j_{+}$  value ( $j_{+} = l_{+} = 1/2$ ) of the single particle final state for a given l value. Namely, it was shown that the  $b_{k}$  coefficients, in the absence of spin distortions, should obey the relationship

$$b_{L}(j_{l})/b_{L}(j_{l}) = -(l+1)/l$$
(1)

and this conclusion (hereafter called j-effect) was found to be consistent with the signs of the b\_2 coefficients known from (p ,  $\gamma$  ) data at isolated energies where E1-capture is dominant.

In the present work the direct-semidirect (DSD) model is used to investigate the *j*-effect for all k and in the whole energy region of the giant multipole resonances. Here some results for radiative capture of 5-35 MeV neutrons by  $^{40}$ Ca, are presented. Calculations are performed<sup>2)</sup> taking into account the IS and IV GQR's and the ISGOR. The single particle final states (and binding energies) considered are : the ground state f 7/2 (8.36 MeV), the first excited state p 3/2 (6.42 MeV) and the *j*\_states

p 1/2 (4.42 MeV), f 5/2 (3.48 MeV). In Fig. 1 the function

$$B_{sum} (E, \theta) = \sum_{k=1}^{6} b_{k} (E) P_{k}^{1} (\cos \theta)$$
(2)

is plotted for capture to the f, states at 11 MeV (upper) and 15 MeV (lower) neutron incident energy and compared with the experimental points<sup>3)</sup> for capture to the f 7/2 state. The opposite signs of the calculated functions  $B_{sum}$  are consistent with eq.(1). Due to the difference in the binding energies of the levels considered,  $E_n = 11 \text{ MeV}$  is near the ISGQR peak for capture to the f 7/2 state and below it for f 5/2. Where E1 radiation dominates, the sum (2) reduces to its second term and the corresponding





Fig. 2. B for the f 7/2 state

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Fig. 3. B<sub>even</sub> for the f 5/2 state

Fig. 4. B<sub>odd</sub> for the p 3/2 state

curves are nearly symmetric. The asymmetry reveals the presence of E2 radiation. Qualitatively similar curves are obtained in the region of the IVGQR (25-30 MeV) as can be seen from Fig. 2. It should be noted however, that at ~ 25 MeV the positive values for  $B_{sum}$  (f 7/2) are pushed toward smaller angles. Similar results are obtained for capture to the  $p_{+}$  states. Therefore, for angles between ~ 30° and 90° the model predicts positive values of  $B_{sum}$  both for capture to  $j_{+}$  and  $j_{-}$  states.

The b-coefficients contain different multipole strength and this limits the possi= bility of using  $B_{sum}$  for studying the *j*-effect. Therefore two functions, which can be related to data from measurements at two supplemental angles, are here introduced :

$$B_{\text{even}}(E,\theta) = \sum_{n} b_{n}(E) P_{n}^{1}(\cos\theta) / [1 + \sum_{n} a_{n}(E) P_{n}(\cos\theta)]$$
(3)

$$B_{\text{odd}}(E,\theta) = \sum_{m} b_{m}(E) P_{m}^{1}(\cos\theta) / [1 + \sum_{n} a_{n}(E) P_{n}(\cos\theta)]$$
(4)

with n = 2,4,6 and m = 1,3,5. Calculations predict positive and negative values of B for capture to  $j_{+}$  and  $j_{+}$  states respectively (see Fig. 3 for f 5/2). This rule holds in the whole energy-angle region for the 4 levels considered, except for the p 3/2 state where, at high energies, near zero-positive values of B even, are obtained.

The rule mentioned agrees with the sign, though values calculated here of the j-ef=fect are greater than those of eq. (1). Figure 4 (p 3/2 state) shows greater varia=tion from zero values at the peaks of IS and IV GQR's thus clearly indicating their positions. Taking into account the difference in binding energies the corresponding  $j_+$  surfaces can be considered symmetric with respect to the zero plane. Calculated

 $J_{\pm}$  surfaces can be considered symmetric with respect to the zero plane. Calculated  $B_{odd}$  agree with the sign and indicate values of the *j*-effect close to those of eq.(1).

The present investigation shows that : 1) the energy-variation of the *j*-effect is closely related to the relative strength of different multipole contributions at a given energy; 2) the predictions of the DSD model together with measured values for the *j*-effect, can be useful as a tool both in assigning *j*-values to final states and in obtaining information about the strength and position of GMR's.

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